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Phase constitution and microstructure of the NbTiVZr refractory high-entropy alloy solidified upon different processing

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ABSTRACT

Time- and temperature-resolved phase formation in the equiatomic NbTiVZr refractory high-entropy alloy has been studied *in situ* using high-energy synchrotron X-ray diffraction and high-speed video imaging during non-equilibrium solidification. Phase formation is shown to be dependent on the solidification conditions. When the melt is undercooled over 80 K it crystallizes as a *bcc* single-phase solid solution despite solute partitioning between dendrites and interdendritic regions. When the sample is for some time kept in the semisolid state, two additional *bcc* phases form in the interdendritic regions. The crystal growth velocity for the NbTiVZr alloy, as estimated from the high-speed videos, shows pronounced sluggish kinetics.

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the strength of this alloy reached a maximum of 500 MPa and then stabilized at 200 MPa during compression testing at 1273 K. On the

other hand, low yield strengths of 75 MPa and 58 MPa at 1273 K

were reported in Refs. [3] and [5], respectively. This perhaps limits its use for high-temperature applications, but it does not exclude

service at moderate temperatures [6]. King et al. [6] used a high

throughput computational tool, Alloy Search and Predict, to screen

over one million four-component equiatomic compositions, looking

for those which would form a single-phase solid solution (SPSS)

at the melting point and have a low thermal neutron absorption

cross-section. The NbTiVZr alloy was found to comply with both

ferent phases and microstructures, depending on the fabrication

and processing route. Senkov et al. [2] produced NbTiVZr samples

by arc-melting followed by hot isostatic pressing at 1473 K and

207 MPa for 2 hours. In sequence, the alloy was homogenized at

1473 K for 24 hours and cooled with a low rate of 10 K·min⁻¹. The microstructure was composed of large grains (\sim 600 μ m) with a body-centered cubic (*bcc*) crystal structure matrix and homogeneously distributed Zr-rich regions of 20 – 30 μ m inside the

grains and sub-micron V-rich particles inside the Zr-rich regions.

The available publications on the NbTiVZr RHEA report on dif-

1. Introduction

The search for materials suitable for high-temperature and radiation-resistant applications gained a new perspective with the introduction of refractory high-entropy alloys (RHEAs) [1]. This design concept opens new avenues for exploring a compositional space based on the high melting temperature metallic elements (W, Ta, Mo, Nb, Hf, V, Zr). At the same time, it brings challenges in investigations of the phase formation and stability ranges, mechanical properties, and corrosion resistance under such extreme thermal conditions. Resolving these key material properties is required to cross the bridge between scientific enthusiasm and practical applications.

The NbTiVZr equiatomic composition was proposed as a low density ($\rho = 6.52 \text{ g}\cdot\text{cm}^{-3}$) RHEA [2,3], showing good compressive ductility and strong work hardening behavior at room temperature. The compression yield strength data for this alloy at the high-temperature regime are scattered. Jia et al. [4] reported that

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requirements.







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Table 1

Measured chemical composition of the NbTiVZr as-cast alloy.

Concentration (at. %)	

The crystal structure of the precipitates could not be identified with X-ray diffraction (XDR) due to their small volume fraction. Yurchenko et al. [7] had also found a small fraction of a second phase at the grain boundaries in NbTiVZr samples homogenized at the same conditions and and observed a splitting of the (110) *bcc* peak on the measured X-ray diffraction (XRD) curves.

In contrast, only one *bcc* phase was found in a NbTiVZr sample prepared by arc-melting of pure components in the work of Tong et al. [8] as well as in a homogenized (1473 K, 100 hours) and water-quenched alloy in the study of King et al. [6]. However, after a subsequent aging at 973 K for 100 hours, the *bcc* SPSS decomposed into a *bcc* matrix and hexagonal closed-packed (*hcp*) and V₂Zr (*C15*-Laves) minor phases [6].

Senkov et al. [5] reported on the appearance of three *bcc* phases after compression deformation of the NbTiVZr alloy with initial SPSS *bcc* structure at 1273 K. It was suggested that there exists a miscibility gap in the NbTiZrV composition, which results in the decomposition of the high-temperature *bcc* phase below 1673 K.

As it has been demonstrated for many HEAs [9–12], processing conditions and kinetics play an important role in determining the final microstructures. A challenge in the study of refractory materials is that high-temperature measurements are usually difficult due to instrumentation limits. For example, measurement of the liquidus (T_1) and solidus (T_s) temperatures cannot be done using conventional differential scanning calorimetry. On the other hand, sample levitation techniques, such as electromagnetic levitation (EML) used in the present work, allow for processing and investigations even in the liquid state.

This study aims at elucidating phase formation and microstructure development upon solidification and during heating/cooling of the NbTiVZr RHEA. A better understanding of metastability, temperature-dependent phase transformations, and the role of kinetics is required for a proper choice of the processing parameters during fabrication and potential use of this material in hightemperature and/or radiation-resistant applications.

2. Materials and methods

Spherical-like samples of the NbTiVZr alloy (about 0.55 - 1 g weight) were produced from pure elements (Ti 99.99 %; Nb, Zr and V 99.8 %) by arc-melting in a water-cooled copper hearth. The furnace was previously evacuated to $5 \cdot 10^{-5}$ mbar and backfilled with Ar up to a pressure of 600 mbar. At least four melting cycles were repeated to achieve chemical homogeneity. A high-purity Ti button (about 10 g weight) placed in the furnace and heated by electric arc was used as an oxygen getter.

The chemical composition of an as-cast sample determined by inductively coupled plasma optical emission spectrometry (ICP-OES) using an iCAP 6500 Duo View from Thermo Fischer Scientific is given in Table 1. The element concentrations deviated by 0.1 to 0.6 at. % from the nominal composition. It should be mentioned that 0.3 at. % of hafnium was also detected in the specimens. The origin of this impurity should be the raw Zr rod. Hf and Zr have the same crystal structure, comparable electronegativity, atomic radius,

and similar physical properties [13], which leads us to believe that the small amount of Hf will not influence the results of this study.

The mobile EML facility of the IFW Dresden [14] was used to carry out the solidification experiments. The samples were levitated and inductively heated using a specially designed water-cooled coil made of copper tube with an inner diameter of 1 mm and an outer diameter of 2 mm. A high-frequency generator TruHeat HF 3010 from Hüttinger Elektronik (10 kW) was operated at 280 kHz. The experimental chamber was evacuated to better than $5 \cdot 10^{-6}$ mbar and backfilled with high purity He (6N) to 500-700 mbar. The sample cooling was done by two He gas jets (6N purity) directed to the sample surface from opposite sides with a flow of up to 28 l·min⁻¹. A one-color pyrometer, 50 Hz sampling rate, was used for recording temperature-time profiles during sample processing.

In situ synchrotron XRD experiments were carried out at the broadband diffraction beamline P21.1 at the PETRA III synchrotron source at the German Electron Synchrotron DESY, Hamburg. A hybrid pixel PILATUS3 X 2M CdTe detector (DECTRIS) at a sampling rate of 20 Hz was used to acquire the diffracted intensity in transmission mode. The recorded 2D images were azimuthally integrated with the FIT2D software [15], to obtain XRD intensities as functions of the diffraction vector $Q = 4\pi \sin \theta / \lambda$, where 2θ is the diffraction angle and λ is the wavelength (0.122 Å in the present work).

Estimation of dendrite growth velocities v was done by measuring the recalescence front travel distances observed from the *in situ* high-speed videos (HSV) (FASTCAM SA5, 775K-M3, from Photron), measured at 1000 frames per second during crystallization. Traveled distances in pixels were converted to length by using the conversion factor of 68 px·mm⁻¹, which was determined using a standard sample during unit calibration. The videos were analyzed with the Fiji ImageJ software [16]. More details on data analysis can be found in Ref. [17].

Metallographic specimens were prepared using standard grinding techniques and fine polishing with colloidal silica (0.05 μ m) suspension. The microstructural examination was performed using a ZEISS Gemini field emission scanning electron microscope (SEM) equipped with a backscattered electron (BSE) detector and an energy-dispersive X-ray (EDX) detector.

Transmission electron microscopy (TEM) samples were prepared as follows. Firstly, slices cut out from the middle of the samples processed in the EML were thinned by grinding down to ~ 100 µm. Then, thin cross-sections in the range of approximately 50-80 nm (electron-transparent) were obtained from the ground samples by ion beam milling using a focussed ion beam (FIB) in the SEM. The thin cross-sections were fixed onto a TEM grid and placed in a TEM single-tilt holder. The measurements were carried out using a FEI F30 TEM operating in a conventional mode at an accelerating voltage of 300 kV and a Thermo Fisher Scientific Talos F200X operating in conventional and scanning modes at an accelerating voltage of 200 kV. The TEM images were analyzed and processed using the FEI TEM Imaging and Analysis, Thermo Fisher Scientific Velox, and the Gatan Digital Micrograph software. The selected area electron diffraction (SAED) patterns were indexed using the program ELDISCA C# [18].

Room temperature XRD was measured, on the slices cut from the solidified samples, at the P21.1 beamline at DESY (beam cross-section: $1 \times 1 \text{ mm}^2$, $\lambda = 0.122 \text{ Å}$). The diffraction patterns were analyzed using the X'Pert HighScore Plus program [19] and the lattice parameters were calculated with the UnitCell software [20].

Lever rule thermodynamic calculations and Scheil-Gulliver solidification simulation were performed using the Thermo-Calc 2020a software from Thermo-Calc AB [21] and the TCHEA4 database [22]. In both cases, all phases modeled in the TCHEA4 relating to the Nb-Ti-V-Zr system were entered.



Fig. 1. (a) Thermogram and corresponding XRD contour plot measured during heating and subsequent cooling of electromagnetically levitated NbTiVZr alloy; (b) Enlarged part of the XRD contour plot from Fig. 1a showing the crystallization event upon solidification of undercooled melt by $\Delta T = 80$ K; (c) comparison of XRD profiles of solid NbTiVZr alloy at the initial and final stages of EML processing.

3. Results

3.1. In situ synchrotron X-ray diffraction

Fig. 1a shows a temperature-time profile measured during a heating/cooling cycle, coupled with the time-resolved XRD patterns plotted as a color-contour diagram. The pyrometer was calibrated using the melting point of pure vanadium observed on the thermogram measured in the EML facility, at the same processing conditions. Thermal analysis using a calorimeter was not possible since NbTiVZr samples reacted with both Al_2O_3 and ZrO_2 crucibles. The solidus and liquidus temperatures for the NbTiVZr alloy were registered during 16 heating/cooling cycles carried out in the EML facility, and the average temperature and standard deviation were calculated (Table 2). The values of T_s and T_1 from the thermodynamic calculations are also included in Table 2.

In Fig. 1a, a peak appearing at 1766 K on the thermogram measured during heating is due to a change of emissivity upon the onset of melting. The solidus temperature is registered as 1736 K.

Table 2 The solidus T_s and liquidus T_1 temperatures of the NbTiVZr RHEA measured with a one-color pyrometer during heating in the EML facility and calculated using the ThermoCalc AB software and the TCHEA4 database.

	<i>T</i> _s (K)	<i>T</i> ₁ (K)
Experimental CALPHAD	1736 ± 12 1762	$\begin{array}{c} 1858 \pm 7 \\ 1926 \end{array}$

During heating up to T_s only XRD peaks corresponding to a *bcc* crystal structure can be identified. The melting ends at 1858 K and the liquid is further overheated by approximately 300 K above T_1 . During this stage, only diffuse scattering is measured, confirming that the alloy is completely molten. The liquid is then cooled by decreasing the input power to the induction coil and by two jetstreams of He directed to the sample's surface by opposite sides. At 1778 K spontaneous nucleation is recorded by the pyrometer which yields a liquid undercooling of $\Delta T = 80$ K. The XRD peaks that appear following this stage correspond to a bcc crystal structure. Fig. 1c shows the X-ray diffractograms for this particular cycle at two distinct temperatures; during heating at 856 K and cooling at 806 K. It is clear from the patterns that the initial and final states show the same (bcc) structure. A 3D XRD plot corresponding to the crystallization of the sample is shown in Fig. S1 in Supplementary Materials.

Fig. 2a shows a temperature-time profile coupled with the time-resolved XRD patterns for a NbTiVZr sample heated up close to the liquidus temperature, but without reaching a complete melting. This is seen from the (110), (200), and (211) peaks that extend throughout the whole duration of the experiment. During heating from room temperature up to T_s , only diffraction peaks that are indexed to a *bcc* structure are recognizable. In the semisolid state, between about 15 s and 75 s, Bragg peaks of the *bcc* phase overlap with the diffuse peaks scattered by the liquid phase, see for example the XRD pattern at 1830 K and ~56 s in Fig. 2c.

Upon the sample cooling, starting from approximately 75 s, reflections of several crystalline phases are observed in the XRD patterns, see the enlarged plot in Fig. 2b. These peaks are indexed to



Fig. 2. (a) Thermogram and corresponding XRD contour plot measured from a NbTiVZr sample heated between T_1 and T_s and subsequently cooled; (b) enlarged part of the XRD contour plot from Fig 2a showing the appearance of different diffraction peaks upon cooling; (c) comparison of XRD profiles at different stages of EML processing.

two extra *bcc* phases. The onset temperature for the appearance of the *bcc*#2 is 1491 K and the onset temperature for the *bcc*#3 phase is 1281 K. A small peak on the temperature-time profile around 74 s (Fig. 2a) corresponds to a change of the emissivity of the alloy upon crystallization. A 3D XRD plot measured during cooling of this sample is shown in Fig. S2 in supplementary materials.

The samples were always weighed before and after the levitation experiments and the maximum weight loss was \sim 0.01 wt. %. This indicates that neither oxidation nor crucial compositional changes occurred.

3.2. Room temperature synchrotron X-ray diffraction

Fig. 3 displays the *ex situ* synchrotron XRD measurements of the samples shown in Figs. 1 and 2 after solidification. In Fig. 3a only peaks corresponding to a *bcc* SPSS (*bcc*#1) are seen. In Fig. 3b besides the peaks corresponding to the primary *bcc* phase (*bcc*#1), additional (111) reflections are marked for *bcc*#2 and *bcc*#3 phases. Fig. 3c shows the same pattern as in Fig. 3b but with a reduced scale of intensity, so the weaker reflections from other crystallographic planes are distinguishable. The lattice parameters extracted from the XRD curve are 0.330 nm, 0.312 nm and 0.344 nm for the *bcc*#1, *bcc*#2 and *bcc*#3 phase, respectively. These values are in a good agreement with those of the three *bcc* phases found in the NbTiVZr alloy after high-temperature compression test by Senkov et al. [5]: 0.332 nm, 0.324 nm and 0.345 nm.

3.3. High-speed video imaging

Digital images representing the solidification of a NbTiVZr melt at an undercooling of 107 K are shown in Fig. 4a. In the upper left panel, the melt at its maximum undercooling is seen. In the uppermiddle panel, a phase can be observed (brighter contrast) growing in the liquid (darker contrast). The remaining panels show the sequence of the crystallization process. The difference in the contrast is due to the release of latent heat from the crystallizing solid. In Fig. 4b the variation of the images' grayscale level, measured on a 50×50 pixel area in the central part of a sample in the videos, is shown as a function of time during the crystallization of the liquid. The grayscale dataset is smoothened because the crystallization is slower than the rotation of the sample. This makes either the liquid or solid phase to be within the analyzed area in a certain period of time. Therefore, a periodic pattern of the grayscale is observed especially once the crystallization starts. The video is available in Supplementary Materials (Video 1).

3.4. Microstructure and structure analysis

SEM micrographs recorded using the BSE detector, combined with the corresponding EDX elemental maps, are displayed in Figs. 5a and 5b for the sample shown in Fig. 3a. Fig. 5a shows a dendritic microstructure with solute rejection to interdendritic regions (brighter contrast). From the EDX elemental mappings, it is possible to observe that the dendrite cores are rich in Nb and



Fig. 3. Room temperature XRD patterns of the NbTiVZr alloy measured with high-energy synchrotron radiation: a) sample solidified from the liquid state with a melt undercooling $\Delta T = 80$ K; b) sample solidified from the semi-solid state; c) XRD pattern from panel (b) with a reduced Intensity-scale to highlight the reflections of the three *bcc* phases.

Ti, while they are depleted of Zr and V. Despite the evident microsegregation (Fig. 5b), no second phase can be recognized from the SEM image (Fig. 5a) which is in agreement with the *in situ* (Fig. 1) and *ex situ* synchrotron XRD measurements in Fig. 3a.

In Fig. 5c, a dendritic microstructure, which is much coarser than that shown in Fig. 5a, is observed with three different contrast regions from the SEM-BSE image: i) a medium-bright region corresponding to the matrix, ii) a brighter region along the grain boundaries, and iii) submicron-sized dark particles inside the latter. The SEM-EDX mappings show that the matrix, *bcc*#1 in the XRD profile in Fig. 3b, is rich in Nb and Ti, while depleted of Zr and V. The bright regions along the grain boundaries seen in the SEM-BSE image (Fig. 5c) are enriched with Zr and V, and the dark particles are rich with vanadium. SEM and EDX images with higher magnification are shown in Fig. S3 in Supplementary Material. The crystal structure of the precipitates was further investigated by TEM.

Fig. 6a shows a SEM-BSE image of the dendritic microstructure of a sample solidified with a melt undercooling of 80 K, as illustrated in Fig. 5a. The red line marks the area from which a FIB-SEM cross-section for further TEM investigations was prepared. A low magnification TEM image of the prepared cross-section is shown in Fig. 6b. Multiple SAED measurements were acquired to confirm the presence of a single *bcc* phase. Typical electron diffractograms for different regions of the sample are shown in Fig. 6c and Fig. 6d. The diffraction patterns were indexed as a *bcc* phase with a lattice parameter of 0.330 nm corresponding to the lattice parameter determined from the synchrotron XRD at room temperature (Fig. 3a). All patterns were indexed corresponding to the (111) zone axis. Fig. 6e shows images of the cross-sectional specimens collected with a high-angle annular dark-field (HAADF) detector with the TEM operating in the scanning mode (STEM). STEM is a variation of TEM that uses a focused electron beam that is scanned over a thin specimen. STEM-HAADF imaging allows the differentiation between different atomic elements as in this imaging mode, the intensity scales with approximately the atomic number squared ($I \sim$ Z^2). STEM also makes it possible to probe the local chemical composition of the specimen using the EDX detector. STEM-EDX mappings of different regions of the sample (Fig. 6f and Fig. 6g) show the homogeneous nature of the sample. A thorough investigation of the cross-section was carried out with high-resolution images and multiple regions for SAED were acquired and are presented in Supplementary Material (Figs. S4 and S5).

Fig. 7 shows images of a sample processed from the semisolid state, as illustrated in the XRD pattern in Fig. 3b. The SEM-BSE image is shown in Fig. 7a. To investigate these three regions in detail, a cross-section indicated by the red line in Fig. 7a has been studied by TEM. Low magnification and high magnification TEM images depicting the different regions of interest are shown in Fig. 7b and Fig. 7c, respectively. Figs. 7d and 7e show electron diffraction patterns corresponding to the *bcc* matrix (lattice parameter: 0.330 nm). A thorough study of the matrix was also conducted in this



Fig. 4. a) Digital images representing the solidification of a NbTiVZr melt with an undercooling $\Delta T = 107$ K; the time scale was set to zero at a moment shortly before the onset of crystallization; b) variation of the grayscale level on a 50 × 50 pixel area as a function of time measured in the central part of the sample shown in panel (a).

case with a high-resolution image and multiple SAED patterns and is presented in Supplementary Material (Fig. S6). Fig. 7f shows a diffraction pattern on the interface between the Zr-rich region (lattice parameter: 0.344 nm) and the V-rich region (lattice parameter: 0.312 nm), where it is possible to distinguish two different diffraction points. Fig. 7g presents the diffraction pattern of the Vrich particles (lattice parameter: 0.312 nm). All patterns were indexed corresponding to the (100) zone axis of a *bcc* crystal structure, showing that all grains and interfaces are aligned in the same *bcc* superstructure. The EDX elemental mapping of the different regions depicted in the HAADF image (Fig. 7h) shows that Nb, Zr, Ti, and V are homogeneously distributed in the matrix (Fig. 7i), whereas Fig. 7j clearly shows the V-rich and Zr-rich regions at the grain boundaries. This is in agreement with the EDX maps obtained with SEM (Fig. 5d). High-resolution images of all phases are presented in the Supplementary Materials (Fig. S7), while the elemental mapping images of higher magnifications are presented in Fig. S8.

4. Discussion

4.1. Phase formation upon solidification

To the best of our knowledge, no experimental values of T_s and T_1 have been reported in the literature for this RHEA. This is supposedly due to the reactivity of this alloy with the crucible materials during DTA experiments at high temperatures, as observed in the present work. The solidus temperature $T_s = 1736 \pm 12$ K



Fig. 5. Microstructural characterization of the NbTiVZr solidified samples: a, b) SEM-BSE image and corresponding element maps measured from the sample solidified with a melt undercooling $\Delta T = 80$ K; c, d) SEM-BSE image and elemental maps measured from the sample solidified from the semisolid state.

(Fig. 1 and Table 2) as determined using the pyrometer in the EML device, is in good agreement (only 26 K lower) with the value predicted for equilibrium conditions by CALPHAD, as shown in Fig. 8 and Table 2. The liquidus temperature $T_1 = 1858 \pm 7$ K is 68 K lower than that predicted by CALPHAD (Fig. 8 and Table 2). Usually T_s is much less dependent on heating rates and sample sizes than T_1 [23]. Considering the mean value of T_s and the standard deviation in Table 2, one can see a very good agreement with CALPHAD prediction. On the other hand, T_1 should be much more dependent on the heating rate and sample mass [23], but even so, the values are in reasonable agreement.

The *in situ* XRD patterns measured during crystallisation of the NbTiVZr liquid undercooled by 80 K in the present work (Fig. 1) show only peaks of one *bcc* phase (*bcc#1*). The microstructural and structural analysis (Figs. 5a and 6) further confirm a homogenous nature of the SPSS despite the microsegregation of Zr and V observed in the interdendritic regions (Fig. 5b).

The sequence of images in Fig. 4a reveals that multiple grains nucleate and grow in the undercooled melt simultaneously and that they collide as solidification progresses. It is obvious from the plot of the gray value versus time (Fig. 4b) that only one recalescence (quasi-adiabatic release of heat upon crystallization) occurs. Thus, considering the *in situ* XRD showed in Fig. 1, the crystallizing solid corresponds to the dendrites of a single *bcc#1* phase.

It is safe to conclude that the only transformation that occurred was from the liquid into a *bcc#1* SPPS when the melt was undercooled by 80 K. This agrees with previous studies where a *bcc* SPSS was obtained by quenching of NbTiVZr alloy from the liquid state [8] or homogenized state [6]. A wide stability range of the primary *bcc#1* phase below T_1 also follows from the thermodynamic calculations (Fig. 8). Formation of a second *bcc#2* phase at equilib

rium conditions is predicted by CALPHAD calculations to occur *via* a spinodal decomposition, starting at 1041 K.

When the NbTiVZr sample was cooled from the semisolid state, formation of three bcc phases were observed by in situ XRD measurements (Fig. 2). The two bcc phases were revealed at the grain boundaries of the primary bcc phase by the microstructure analvsis of the soldified samples (Figs. 3b, 5c, and 7). From the onset temperatures of the bcc#2 and bcc#3 (Fig. 2b), which are respectively 245 K and 455 K below T_s , one could conclude that they occurred as solid-state transformations. However, during nonequilibrium solidification constitutional undercooling may lead to an overall change in the composition of the remaining liquid. The interdendritic liquid then solidifies at a different T_s than the nominal alloy's composition. Furthermore, the pyrometer only measures a small surface area on the north pole of the droplet. It cannot distinguish temperature gradients over the sample and it does not probe local temperatures. The formation of bcc#2 and bcc#3 phases (Fig. 2) occurred during the final stage of solidification with crystallization of the interdendritic liquid. This is evident from the solidified microstructures shown in Figs. 5c and 5d. While the sample was cooled between T_1 and T_s solute partitioning from the primary dendrites (bcc#1) to the interdendritic regions created the compositional conditions for their formation.

The *ex situ* XRD patterns (Figs. 3a and 3b) confirm that in both cases the phases formed at high temperature, as shown from the *in situ* XRD patterns in Figs. 1 and 2, are preserved at room temperature. The formation of equilibrium *hcp* and *C15*-Laves phases which are predicted by CALPHAD (see Fig. 8) and have been experimentally observed after long-time annealing, at lower temperatures [6,7], were kinetically suppressed due to the fast cooling rate (~100 K·s⁻¹) in the solid-state.



Fig. 6. a) SEM-BSE image showing the dendritic microstructure of the NbTiVZr sample solidified with a melt undercooling $\Delta T = 80$ K; the red line marks the area from which the TEM cross-section was prepared; b) TEM image of the specimen cross-section; c, d) SAED patterns for two different regions of the cross-section; e) STEM-HAADF image of the cross-section; f, g) the corresponding STEM-EDX elemental maps for the different regions shown in panel (e) (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article).



Fig. 7. a) SEM image showing different phases in the NbTiVZr alloy solidified from the semisolid state: the dark contrast corresponds to the V-rich region (lattice parameter: 0.312 nm), the brighter contrast corresponds to the Zr-rich region (lattice parameter: 0.344 nm), and the medium-bright contrast corresponds to the matrix (lattice parameter: 0.330 nm). The red line marks the area from which a TEM cross-section was prepared; b) TEM image of the cross-section; c) high-magnification TEM image showing a region with all phases coexisting; d-g) SAED patterns for the matrix, the interface between V-rich and Zr-rich region, and the V-rich region, respectively; h) STEM-HAADF image of the cross-section; i, j) local elemental maps for different regions. The elemental mapping of panel i (matrix) is homogeneous, whereas the maps in panel j show the different chemical compositions (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article).

In the quaternary Nb-Ti-V-Zr system, three of the the six constituent binary systems (Ti-V, Nb-Zr and Nb-V) are characterized by a solid-state miscibility gap with a critical temperature in the vicinity of the equiatomic composition. Below the critical temperature, the primary *bcc* phase decomposes into Ti-rich *bcc* and V-rich *bcc* (Ti-V system), Nb-rich *bcc* and Zr-rich *bcc* (Nb-Zr system), and Nb-rich *bcc* and V-rich *bcc* (Nb-V system). These constituent binaries contribute to the quaternary system such that theoretically at least three *bcc* composition sets should exist: Nbrich (*bcc*#1), V-rich (*bcc*#2) and (Zr,Ti)-rich *bcc*#3 (Zr- and Ti-rich *bcc* combines due to unlimited solid solubility of Zr in Ti and vice versa).

Non-equilibrium solidification of the NbTiVZr alloy calculated using the Scheil-Gulliver model is shown in Fig. 9. The solidification proceeds between $T_1 = 1926$ K, which is the same as for the equilibrium conditions, and a eutectic temperature $T_e \approx 1491$ K at what the fraction of the liquid phase decreases to about 1 %. The temperature dependences for the compositions of the liquid and



Fig. 8. Phase fraction as a function of temperature for equilibrium solidification of the NbTiVZr alloy predicted using the ThermoCalc AB software and the TCHEA4 database.



Fig. 9. The temperature dependence of the fraction of solid phases during solidification of the NbTiVZr alloy under the equilibrium (lever rule) and non-equilibrium (Scheil-Gulliver) conditions simulated with the ThermoCalc AB software using the TCHEA4 database.

the solid phases are shown in Fig. 10. The Scheil-Gulliver simulation suggests that cores of the *bcc#1* dendrites formed in earlier stage of solidification are enriched with Nb and Ti (Fig. 10b, *T* above ~1770 K). Upon cooling, both V and Zr continuously segregate into the interdendritic liquid (Fig. 10a) and start to dominate in the peripheral areas of the dendrites (Fig. 10b, *T* below ~1735 K). Below 1492 K, the interdendritic liquid solidifies into the V-rich *bcc#2* and a Zr-rich *bcc#3* phases *via* eutectic reaction which completely finishes at 1491 K. These results agree with the element distribution maps presented in Fig. 5d and in Fig. S3 in Supplementary Material. Also, such solidification path conforms with

the in-situ XRD measurement upon solidification of the NbTiVZr alloy from the semisolid state (Fig. 2). In that case, the (110) reflections of a V-rich phase was observed to appear at 1491 K. The (110) peak of a Zr-rich phase was noticed at a lower temperature (1281 K). The discrepancy on the formation of the metastable *bcc* phases between the Scheil-Gulliver calculation and in-situ XRD measurements could be explained either by some deviation of the real and simulated solidification paths or by somewhat imperfect modelling of the TCHEA4 database (see for this also our previous papers [17,24]), or by experimental uncertainty. Indeed, the molar fraction of the *bcc*#2 and *bcc*#3 phases is very small and they



Fig. 10. Composition of the liquid and solid phases predicted by Scheil-Gulliver solidification simulation of NbTiVZr.

Table 3

Crystal growth velocity v as a function of melt undercooling ΔT upon crystallization of the NbTiVZr RHEA.

ΔT (K)	$v (m \cdot s^{-1})$
77	~0.009
101	~0.015
107	~0.017
112	~0.019
145	~0.037

can be noticed on XRD patterns only when they reach a detectable quantity at the X-ray beam position.

4.2. Crystal growth kinetics

It can be seen from the high-speed video of the crystallizing NbTiVZr melt at 107 K undercooling (Video 1 in Supplementary Materials), as well as from the selected digital images in Fig. 4 that the crystal growth velocity v is extremely low. The values of v extracted from the digital videos taken after different melt undercooling are given in Table 1. The largest v-value, measured at 145 K, is equal just to about 0.04 m·s⁻¹. Table 3

It is known that the growth velocity of *bcc* crystals is in general smaller than that of *fcc* crystals. For example it has been shown in a recent study of Andreoli *et al.* [17] that v_{bcc} is about 3 – 4 times smaller than v_{fcc} in the CoCrFeNi HEA at ~80 K undercooling, relative to the liquidus temperature of the respective phase. However, the growth velocity of the *bcc* phase in the undercooled CoCrFeNi

melt is about two orders of magnitude larger than in the NbTiVZr melt. For compatison, $v_{\rm bcc}$ is equal to ~1.15 m·s⁻¹ at $\Delta T = 80$ K in CoCrFeNi and ~0.01 m·s⁻¹ at $\Delta T = 77$ K in NbTiVZr.

Wang et al. [25] have recently studied the solidification of electrostatically levitated NbMoTaWZr refractory HEA and reported a growth velocity for the *bcc* phase of 0.23 m·s⁻¹ for $\Delta T = 82$ K. This is also very low growth rate. Yet, it is about one order of magnitude larger than that obtained for the NbTiVZr alloy at a similar undercooling in the present study.

Probably the difference in the growh velocities of the *bcc* phases in NbTiVZr and NbMoTaWZr can be explained to some extent by the different methods used for their determination. In Ref. [25] the growth velocity was measured by the infrared method as the ratio of the sample diameter to the recalescence time. In the present work, it is calculated from the crystallisation front propagation over time using high-speed video. It is, however, quite plausible that the crystal growth kinetics of the NbTiVZr alloy is indeed sluggisher than that of the NbMoTaWZr, at least at relatively low undercooling, as in the present study, where the solidification is diffusion-controlled. For example, it has been shown by Vaidya *et al.* [26,27] that the diffusion rate in HEAs, in the solid-state, does not necessarily decrease with increasing number of alloy constituents N as it has been supposed in the first publications on HEAs.

4.3. Microstructural features

At a very low undercooling, which must be the case when the melt was slowly cooled from the semisolid state (Fig. 2), constitu-

tional undercooling is dominant. At this level of undercooling the dendrite tip is capable of coarsening which can lead to coarser microstructures. This is observed when comparing the dendritic structures shown in Figs. 5a and 5c. Crystal growth velocity should be very limited which also contributes to coarsening of the microstructure. Unfortunately, the visualization of the solidification when samples are processed in this manner cannot be done because of the extremely weak contrast between the liquid and the crystalline phase due to a small release of latent heat. Under these conditions, solute rejection occurring at the solid/liquid interface has had sufficient time to result in partitioning which is evident from the interdendritic regions in Figs. 5c and 5d. Contrarily, the microstructure of the sample solidified under an undercooling of 80 K revealed a much finer dendritic structure and less microsegregation (Figs. 5a and 5b). At a moderate undercooling, such as $\Delta T = 80$ K, the decrease of dendrite tip and faster crystal growth kinetics lead to smaller dendrites (Fig. 5a).

5. Conclusions

The initial state and solidification kinetics plays a central role in the high-temperature phase formation of the NbTiVZr RHEA. While a *bcc* single-phase microstructure was formed when the liquid was undercooled over 80 K, solidification from the semisolid state between T_1 and T_s lead to the development of a multi-phase microstructure. In the first case, as the driving force for crystallization is proportional to the liquid undercooling, a non-equilibrium solidification is established with (relatively) higher crystal growth rates which were evident from the fine solidified microstructure. In the latter case, with negligible undercooling, a coarse dendritic microstructure was formed with intense microsegregation of vanadium and zirconium into the interdendritic regions and appearance of two additional *bcc* phases at the grain boundaries.

The solute microsegregation and formation of the *bcc*#2 and *bcc*#3 phases could be explained by Scheil-Gulliver solidification that was simulated by the CALPHAD approach using the TCHEA4 database. The discrepancy on the onset temperature for the formation of the *bcc*#3 phase between CALPHAD prediction and experiment can be explained either by deviation of the real and simulated solidification paths, somewhat imperfect modelling of the TCHEA4 database or experimental uncertainty.

The growth velocity of the primary *bcc* phase in the NbTiVZr alloy is found to be about one order of magnitude smaller than that of the NbMoTaWZr alloy at a similar undercooling reported in the literature.

Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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Supplementary materials

Supplementary material associated with this article can be found, in the online version, at doi:10.1016/j.actamat.2021.117416.

References

- O.N. Senkov, D.B. Miracle, K.J. Chaput, J.-P. Couzinie, Development and exploration of refractory high entropy alloys—A review, J. Mater. Res. 33 (2018) 1–37, doi:10.1557/jmr.2018.153.
- [2] O.N. Senkov, S.V. Senkova, C. Woodward, D.B. Miracle, Low-density, refractory multi-principal element alloys of the Cr-Nb-Ti-V-Zr system: Microstructure and phase analysis, Acta Mater 61 (2013) 1545–1557, doi:10.1016/j.actamat. 2012.11.032.
- [3] O.N. Senkov, S. Rao, K.J. Chaput, C. Woodward, Compositional effect on microstructure and properties of NbTiZr-based complex concentrated alloys, Acta Mater 151 (2018) 201–215, doi:10.1016/j.actamat.2018.03.065.
- [4] Y. Jia, L. Zhang, P. Li, X. Ma, L. Xu, S. Wu, Y. Jia, G. Wang, Microstructure and Mechanical Properties of Nb-Ti-V-Zr Refractory Medium-Entropy Alloys, Front. Mater. 7 (2020) 1–11, doi:10.3389/fmats.2020.00172.
- [5] O.N. Senkov, S.V. Senkova, D.B. Miracle, C. Woodward, Mechanical properties of low-density, refractory multi-principal element alloys of the Cr-Nb-Ti-V-Zr system, Mater. Sci. Eng. A. 565 (2013) 51–62, doi:10.1016/j.msea.2012.12.018.
- [6] D.J.M. King, S.T.Y. Cheung, S.A. Humphry-Baker, C. Parkin, A. Couet, M.B. Cortie, G.R. Lumpkin, S.C. Middleburgh, A.J. Knowles, High temperature, low neutron cross-section high-entropy alloys in the Nb-Ti-V-Zr system, Acta Mater 166 (2019) 435–446, doi:10.1016/j.actamat.2019.01.006.
- [7] N. Yurchenko, N. Stepanov, M. Tikhonovsky, G. Salishchev, Phase Evolution of the AlxNbTiVZr (x = 0; 0.5; 1; 1.5) High Entropy Alloys, Metals (Basel) 6 (2016) 298, doi:10.3390/met6120298.
- [8] Y. Tong, S. Zhao, H. Bei, T. Egami, Y. Zhang, F. Zhang, Severe local lattice distortion in Zr- and/or Hf-containing refractory multi-principal element alloys, Acta Mater 183 (2020) 172–181, doi:10.1016/j.actamat.2019.11.026.
- [9] F. He, Z. Wang, Q. Wu, J. Li, J. Wang, C.T.T. Liu, Phase separation of metastable CoCrFeNi high entropy alloy at intermediate temperatures, Scr. Mater. 126 (2017) 15–19, doi:10.1016/j.scriptamat.2016.08.008.
- [10] Q. Tang, Y. Huang, H. Cheng, X. Liao, T.G. Langdon, P. Dai, The effect of grain size on the annealing-induced phase transformation in an Al0-3CoCrFeNi high entropy alloy, Mater. Des. 105 (2016) 381–385, doi:10.1016/j.matdes.2016.05. 079.
- [11] G. Laplanche, S. Berglund, C. Reinhart, A. Kostka, F. Fox, E.P. George, Phase stability and kinetics of *σ*-phase precipitation in CrMnFeCoNi high-entropy alloys, Acta Mater 161 (2018) 338–351, doi:10.1016/j.actamat.2018.09.040.
- [12] E.J. Pickering, R. Muñoz-Moreno, H.J. Stone, N.G. Jones, Precipitation in the equiatomic high-entropy alloy CrMnFeCoNi, Scr. Mater. 113 (2016) 106–109, doi:10.1016/j.scriptamat.2015.10.025.
- [13] H. Jiang, R.I. Gomez-Abal, P. Rinke, M. Scheffler, Electronic band structure of zirconia and hafnia polymorphs from the GW perspective, Phys. Rev. B. 81 (2010) 085119, doi:10.1103/PhysRevB.81.085119.
- [14] O. Shuleshova, I. Kaban, W. Loser, S. Ziller, U. Reinhold, D. Lindackers, Metastable phase formation in undercooled Fe-Co melts under terrestrial and microgravity conditions, IFW Dresden, Annu. Rep. (2017) 39–42 https://www. ifw-dresden.de/uploads/users/10/uploads/files/JB_2017.pdf.
- [15] A.P. Hammersley, S.O. Svensson, M. Hanfland, A.N. Fitch, D. Hausermann, D. Häusermann, Two-dimensional detector software: From real detector to idealised image or two-theta scan, High Press. Res. 14 (1996) 235–248, doi:10. 1080/08957959608201408.
- [16] J. Schindelin, I. Arganda-Carreras, E. Frise, V. Kaynig, M. Longair, T. Pietzsch, S. Preibisch, C. Rueden, S. Saalfeld, B. Schmid, J.-Y. Tinevez, D.J. White, V. Hartenstein, K. Eliceiri, P. Tomancak, A. Cardona, Fiji: an open-source platform for biological-image analysis, Nat. Methods. 9 (2012) 676–682, doi:10. 1038/nmeth.2019.
- [17] A.F. Andreoli, O. Shuleshova, V.T. Witusiewicz, Y. Wu, Y. Yang, O. Ivashko, A.-C. Dippel, M.v. Zimmermann, K. Nielsch, I. Kaban, In situ study of nonequilibrium solidification of CoCrFeNi high-entropy alloy and CrFeNi and CoCrNi ternary suballoys, Acta Mater 212 (2021) 116880, doi:10.1016/j.actamat. 2021.116880.
- [18] J. Thomas, T. Gemming, ELDISCA C# a new version of the program for identifying electron diffraction patterns, in: W.T. Luysberg M., Tillmann K. (Ed.), EMC 2008 14th Eur. Microsc. Congr. 1–5 Sept. 2008, Aachen, Ger., Springer Berlin Heidelberg, Berlin, Heidelberg, 2008: pp. 231–232. doi:10.1007/978-3-540-85156-1_116.
- [19] T. Degen, M. Sadki, E. Bron, U. König, G. Nénert, The high score suite, Powder Diffr (2014), doi:10.1017/S0885715614000840.
- [20] T.J.B. Holland, S.A.T. Redfern, UNITCELL: A nonlinear least-squares program for cell-parameter refinement implementing regression and deletion diagnostics, J. Appl. Crystallogr. (1997), doi:10.1107/S0021889896011673.
- [21] Thermo-Calc AB Software, (2020). https://thermocalc.com/.
- [22] H.-L. Chen, H. Mao, Q. Chen, Database development and Calphad calculations for high entropy alloys: Challenges, strategies, and tips, Mater. Chem. Phys. 210 (2018) 279–290, doi:10.1016/j.matchemphys.2017.07.082.

- [23] M. Žaludová, B. Smetana, S. Zlá, J. Dobrovská, K. Gryc, K. Michalek, R. Dudek, Study of DTA method experimental conditions and of their influence on obtained data of metallic systems, in: Met. 2012 Conf. Proceedings, 21st Int. Conf. Metall. Mater, 2012, pp. 640–645.
 [24] O. Stryzhyboroda, V.T. Witusiewicz, S. Gein, D. Röhrens, U. Hecht, Phase Equi-
- [24] O. Stryzhyboroda, V.T. Witusiewicz, S. Gein, D. Röhrens, U. Hecht, Phase Equilibria in the Al-Co-Cr-Fe-Ni High Entropy Alloy System: Thermodynamic Description and Experimental Study, Front. Mater. 7 (2020), doi:10.3389/fmats. 2020.00270.
- [25] W.L. Wang, L. Hu, S.J. Yang, A. Wang, L. Wang, B. Wei, Liquid Supercoolability and Synthesis Kinetics of Quinary Refractory High-entropy Alloy, Sci. Rep. 6 (2016) 1–7, doi:10.1038/srep37191.
- [26] M. Vaidya, K.G. Pradeep, B.S. Murty, G. Wilde, S.V. Divinski, Bulk tracer diffusion in CoCrFeNi and CoCrFeMnNi high entropy alloys, Acta Mater 146 (2018) 211–224, doi:10.1016/j.actamat.2017.12.052.
- [27] M. Vaidya, G.Mohan Muralikrishna, S.V. Divinski, B.S. Murty, Experimental assessment of the thermodynamic factor for diffusion in CoCrFeNi and CoCrFeMnNi high entropy alloys, Scr. Mater. 157 (2018) 81–85, doi:10.1016/j.scriptamat.2018.07.040.